

### Clebsch-Gordan Coefficient Example Phys 401

Consider the total angular momentum operator for the Hydrogen atom:  $\vec{J} = \vec{L} + \vec{S}$ , where  $\vec{L}$  is the orbital angular momentum of the electron and  $\vec{S}$  is the spin angular momentum of the electron. The eigenfunctions of  $J^2$  can be expressed as linear combinations of states with different values of  $m_\ell$  and  $m_s$  using the world-famous Clebsch-Gordan coefficients ( $C_{m_\ell m_s}^{\ell s j m_j}$ ) as:

$$|j m_j\rangle = \sum_{m_\ell+m_s=m_j} C_{m_\ell m_s}^{\ell s j m_j} |\ell m_\ell\rangle |s m_s\rangle \quad [4.183]$$

where the ket  $|\ell m_\ell\rangle$  represents the spherical harmonics  $Y_\ell^{m_\ell}$ . The C-G coefficient values are given in Table 4.8 on page 179 of Griffiths. Remember that all of the coefficients should appear under a square root, with the minus sign (if any) out front.

Now for an example of how to construct states that are simultaneous eigenfunctions of  $L^2$ ,  $S^2$ ,  $J^2$  and  $J_z$ . Take the case again of hydrogen with  $\ell=1$  and spin  $s=1/2$ . How do we find the state with  $j=3/2$  and  $m_j=-1/2$  in terms of the  $Y_\ell^{m_\ell}$  and spinors? Look at the  $1 \times 1/2$  CG Table on page 179. We are led to this table because we are combining an angular momentum vector with  $\ell=1$  and spin vector with  $s=1/2$ .

$1 \times 1/2$	$3/2$				
	$+3/2$	$3/2$	$1/2$		
	$+1$	$+1/2$	$1$	$+1/2$	$+1/2$
		$+1$	$-1/2$	$1/3$	$2/3$
		$0$	$+1/2$	$2/3$	$-1/3$
				$0$	$-1/2$
				$-1$	$+1/2$
				$2/3$	$1/3$
				$1/3$	$-2/3$
				$-1$	$-1/2$
				$3$	$2$
$2 \times 1$	$3$	$3$	$2$	$1$	$1$
	$+3$	$3$	$2$	$1$	$1$

Now look under the column labeled “ $\begin{matrix} 3/2 \\ -1/2 \end{matrix}$ ”. It says:

$$\left| \begin{matrix} 3 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle = \sum_{m_\ell+m_s=-1/2} C_{m_\ell m_s}^{1 \ 1/2 \ 3/2 \ -1/2} |1 m_\ell\rangle \left| \begin{matrix} 1 \\ 2 \end{matrix} m_s \right\rangle$$

$$\left| \begin{matrix} 3 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle = \sqrt{\frac{2}{3}} |1 \ 0\rangle \left| \begin{matrix} 1 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle + \sqrt{\frac{1}{3}} |1 \ -1\rangle \left| \begin{matrix} 1 \\ 2 \end{matrix} \ \frac{1}{2} \right\rangle$$

This can be written in a more familiar way in terms of spherical harmonics and spinors as:

$$\left| \begin{matrix} 3 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle = \sqrt{\frac{2}{3}} Y_1^0 \chi_- + \sqrt{\frac{1}{3}} Y_1^{-1} \chi_+$$

One can move back and forth between the coupled and un-coupled representations using the Clebsch-Gordan table on page 179. Here is the schematic

layout for the CG table for combining two spins (called  $\vec{S}_1, \vec{S}_2$ ) to form a total spin  $\vec{S} = \vec{S}_1 + \vec{S}_2$  ( $S^2$  has eigenvalue  $s(s+1)\hbar^2$ ):

General Schematic of the C-G Table

$$\vec{S} = \vec{S}_1 + \vec{S}_2$$

$$S_1 \times S_2$$

$$\begin{array}{|c|} \hline S \\ \hline m_s \\ \hline \end{array}$$

Coupled  
Representation

$$\begin{array}{|cc|} \hline m_{s_1} & m_{s_2} \\ \hline \end{array}$$

$CG\#$

Un-Coupled  
Representation